

International Journal of Engineering Mathematics: Theory and Application (Online) 1687-6156 <u>http://iejemta.com/</u> VOLUME 4 ISSUE 1

Editorial Team

G. Ahmed

Professor of Computational Engineering Mathematics and Numerical Analysis **Department of Engineering Physics and Mathematics** Associate editor-in-Chief Dr. Hamed Daei Kasmaei PhD in Applied mathematics-Numerical analysis and computational **Department of Mathematics and Statistics**, **Honor President of IEEMS** Mahim Ranian Adhikari **Department of Mathematics Calcutta University** India Carlo Cattani Professor, Tuscia University, Viterbo **Department of Economy and Enterprise DEIM** Italy E-mail: ccattani@unisa.it

Dr. Sunil Kumar National Institute of Technology Jamshedpur Department of Mathematics India Email: skiitbhu28@gmail.com

Praveen Agarwal Ph.D., Professor Anand International College of Engineering Department of Mathematics Jaipur India Email: goyal.praveen2011@gmail.com

Thomas Korimort Mathematician Computer Scientist Dr. tech. Dipl.-Ing AMS University of Leoben Vienna University of Technology Austria Email: tomkori@gmx.net

Dr. Stephen Kirkup

Lecturer in Nuclear Science / Engineering School of Engineering Computing and Technology Building, CM138 University of Central Lancashire United Kingdom Email: smkirkup@uclan.ac.uk

Dr Mehmet Senol Nevsehir Haci Bektas Veli University Department of Mathematics Nev_sehir Turkey Email: msenol@nevsehir.edu.tr

Dr. Muhammad Sadiq Hashmi Associate Professor Department of Computer Science COMSATS Institute of Information Technology Sahiwal Campus Pakistan Email: sadiq.hashmi@gmail.com



Hector Vazquez Leal **Full Time Professor School of Electronic Instrumentation** University of Veracruz Mexico Email: hvazquez@uv.mx Dr. Jvotindra C. Prajapati M.Sc., M. Phil., Ph.D., MIMS, MISTE **Principal, Faculty of Science** Marwadi University **Rajkot-Morbi Highway** RAJKOT- 360003, GUJARAT India Hasan Bulut **Faculty of Science Department of Mathematics Firat University Elazig Turkey** E-mail: hbulut@firat.edu.tr

Fethi Bin Muhammad Belgacem Department of Mathematics Faculty of Basic Education PAAET, Al-Ardhiya Kuwait Email: fbmbelgacem@gmail.com Avishk Mahim Adkhaira Associate Professor of Mathematics Calcutta University India E-mail: math.mra@gmail.com

János Kurdics Professor of Mathematics University of Nyiregyhaza Hungary Academic Member of ATINER Athens E-mail: kurdics@nyf.hu

CONTACT

Professor of Computational Engineering Mathematics and Numerical Analysis Faculty of Engineering Zagazig University Zagazig P. O. 44519 Egypt http://iejemta.com/ Email: sgamil@zu.edu.eg

BEHAVIOR OF THE MAGNITIC FIELD UNDER RANDOM GAZ CLOUD DEFORMATIONS

Shamshiev Fazliddin Tulaevich

National University of Uzbekistan, Department of Astronomy and Astrophysics E-mail: shamshiyev61@gmail.com

https://doi.org/10.5281/zenodo.7313464

Abstract. The clouds of the interstellar medium are considered, which from time to time undergo random deformations due to the influence of the gravitational fields of massive objects or due to uneven braking from the surrounding background. Qualitative behavior of the magnetic field inside of the cloud for a long time t is analyzed.

Keywords: interstellar medium, magnetic field, deformations of a cloud.

Magnetic fields play an important role in many astrophysical processes, such as the movement of charged particles, the evolution of gas clouds and their dynamic evolution [1; 2]. Here we select a model of linear deformations of a homogeneous ellipsoidal cloud. The behavior of the radius of the r_0 vector of a hydrodynamic particle is described by the relation,

$$A_n = A_n A_{n-1} \dots A_1 r_0$$

where r_0 - is the initial value of the radius of the vector, and random matrices A_n describe the cloud deformations in consecutive, also random at the moment t, following each other with an average time interval τ . The theory of multiplication of random matrices predicts that, with probability 1, [3]. The deformable ellipsoid is infinitely stretched into a "needle", i.e., the relations are valid for the semi-axes a_n , b_n , c_n

$$\lim_{n\to\infty}\frac{c_n}{b_n} = \lim_{n\to\infty}\frac{b_n}{a_n} = 0.$$

Moreover, in addition to exceptional degenerate situations, the ratio of half-axes behaves, up to fluctuations, exponentially and there are limits



International Journal of Engineering Mathematics: Theory and Application (Online) **1687-6156** http://iejemta.com/ VOLUME 4 ISSUE 1

$$\lim_{n \to \infty} \frac{\ln a_n}{n} = \alpha, \quad \lim_{n \to \infty} \frac{\ln b_n}{n} = \beta, \quad \lim_{n \to \infty} \frac{\ln c_n}{b_n} = \gamma.$$

If the *abc* is a volume, which is physically natural, is maintained near one value. Then, obviously, $\alpha + \beta + \gamma = 0$. It is possible to imagine that a magnetic field is frozen in the cloud, which is sufficiently weak and does not have a significant reverse effect on the deformation, but is subject to attenuation due to the final conductivity. Let us also assume that the specific conductivity of the surrounding background is significantly less, so that we can safely assume that the field in the environment is zero. The question arises about the qualitative behavior of the magnetic field inside the cloud at a large time t. To do this, it is sufficient to consider a simplified model of such a field, which is almost uniform in the last stages of deformation, when the ellipsoid turns into a thin elliptical disk. The internal field in this case will obviously be directed mainly parallel to the large axis, which can be taken as the x axis, while the z axis is perpendicular to the "pancake" plane. It is convenient to consider the field evolution between the moments of deformation as a one-dimensional process

$$\frac{\partial H_x}{\partial t} = \frac{c^2}{4\pi\sigma} \frac{\partial^2 H_x}{\partial x^2},$$

where σ - is the conductivity, *c* - is the speed of light. The solution of this equation, as is easily verified, goes to the asymptotic mode

$$\frac{\partial H_x}{\partial t} = -\mu H_x, \quad \mu = \frac{c^2}{4\pi\sigma c_0^2}$$

As for the deformations themselves, their smoothed effect would be described by the ratio $H_x \approx \frac{a_n}{a_0} H_0$ or, when switching from discrete time to continuous time for convenience

$$H(t) = k \exp\left(\frac{\alpha t}{\tau}\right), \quad (k = const.) \quad \text{or} \quad \frac{\partial H}{\partial t} = \frac{\alpha}{\tau} H.$$

When the combined effect of dissipation and deformation must take into account both terms in the transition for convenience from discrete time to continuous



International Journal of Engineering Mathematics: Theory and Application (Online) **1687-6156** http://iejemta.com/ VOLUME 4 ISSUE 1

$$\frac{\partial H}{\partial t} = \left(\frac{\alpha}{\tau} - \mu\right) H \approx \left(\frac{\alpha}{\tau} - \frac{c^2}{4\pi\sigma c_0^2}\right) H.$$

Since it is mandatory $\alpha > 0$, and c_0 , on the contrary tends to zero, the value of $\frac{\partial H}{\partial t}$ becomes negative over time. This means that sooner or later the cloud's magnetic field will begin to decrease.

REFERENCES

1. Zeldovich Ya.B., Ruzmaikin A.A., Sokolov D.D., Magnetic fields in astrophysics, New York: Gordon and Breach, 2006.

2. Kadomtsev B. B. Collective phenomena in plasma, M.: 1988.

3. Tutubalin V. N. "On limit theorems for the product of random matrices", Probability Theory and its application., 10:1 (1965), 19-32; Theory Probab. Appl., 10:1 (1965), 15–27.

